

Einstein Field Equation

Let us consider a coordinate system having axes $x^1, x^2, x^3 \dots x^n$. Then we can express the change in some function ϕ_s in terms of these coordinates as

$$d\phi_s = \sum_n \frac{\partial \phi}{\partial x^n} dx^n \quad (1)$$

We can use eqn. 1 with slight adjustment to show how a vector (tensor) transforms. Let a n dimension vector V in the y frame related with the m dimension vector V in the x frame as

$$V_y^n = \sum_m \frac{\partial y^n}{\partial x^m} V_x^m$$

Where m is a dummy variable.

Let us consider combination of m dimensional vector A^m and n dimensional vector B^n . Combination of this vector is a tensor T^{mn} .

$$T^{mn} = A^m B^n \quad (3)$$

In two dimensional space m and n can be either can have value 1 or 2, then T will have 4 versions. For 3 dimensional space, T^{mn} has 9 versions. Using eqn. 2, we can write,

$$T_y^{mn} = \sum_r \frac{\partial y^m}{\partial x^r} A_x^r \sum_s \frac{\partial y^n}{\partial x^s} B_x^s = \sum_{rs} \frac{\partial y^m}{\partial x^r} \frac{\partial y^n}{\partial x^s} A_x^r B_x^s$$

$$T_y^{mn} = \sum_{rs} \frac{\partial y^m}{\partial x^r} \frac{\partial y^n}{\partial x^s} T_x^{rs} \quad (4)$$

This transformation from y to x frame is called contravariant transformation. There is another transformation, covariant transformation which is,

$$T_{mn}^y = \sum_{rs} \frac{\partial x^r}{\partial y^m} \frac{\partial x^s}{\partial y^n} T_{rs}^x \quad (5)$$

According to Pythagoras,

$$ds^2 = dx^{1^2} + dx^{2^2} \dots$$

$$= \sum_{mn} dx^m dx^n \delta_{mn} \quad (5a)$$

Where δ_{mn} is the **Kronecker delta**, is a function of two variables, usually just non-negative integers. The function is 1 if the variables are equal, and 0 otherwise. Now rewriting the equation 1 in slightly different form as

$$dx^m = \sum_r \frac{\partial x^m}{\partial y^r} dy^r = \frac{\partial x^m}{\partial y^r} dy^r$$

Similarly,

$$dx^n = \frac{\partial x^n}{\partial y^s} dy^s$$

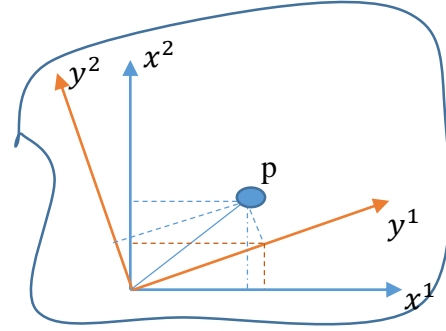


Fig. 1
(2)

Now substituting this value in eqn. 5a,

$$\begin{aligned} ds^2 &= \delta_{mn} \sum_{mn} \frac{\partial x^m}{\partial y^r} dy^r \frac{\partial x^n}{\partial y^s} dy^s \\ &= g_{mn} dy^r dy^s \end{aligned} \quad (6)$$

Where $g_{mn} = \delta_{mn} \sum_{mn} \frac{\partial x^m}{\partial y^r} \frac{\partial x^m}{\partial y^s}$ is called the metric tensor. In flat space g_{mn} simply reduces to Kronecker delta δ_{mn} term. In the Einstein's field equation $g_{\mu\nu}$ terms refers the metric tensor which includes spacetime.

In case of curved surface we have to use the covariant derivatives. Covariant derivative of a tensor is equal to the ordinary derivative plus the correction term (Cristoffel symbol) multiplied by that tensor. Because the tensor $T_{mn}(y)$ equal to covariant derivative of vector V_m . Then

$$T_{mn}(y) = \nabla_n V_m = \frac{\partial V_m}{\partial y^n} + \Gamma_{nm}^r V_r(x) \quad (7)$$

Let us take the covariant derivative of the tensor T_{mn} for two indices from eqn. 7

$$\nabla_p T_{mn} = \frac{\partial T_{mn}}{\partial y^p} + \Gamma_{pm}^r T_{rn} + \Gamma_{pn}^r T_{mr} \quad (8)$$

From the property of tensor we can say that, this covariant derivative of metric tensor will be zero in all frame of reference. Substituting T_{mn} by g_{mn} in equation 8

$$\nabla_p g_{mn} = \frac{\partial g_{mn}}{\partial y^p} + \Gamma_{pm}^r g_{rn} + \Gamma_{pn}^r g_{rm} = 0$$

We can define a tangent vector, the rate of change of distance in a geodesic w.r.to proper time as $\frac{dx^\mu}{d\tau}$. In order for this distance is to be shortest, this derivative should be zero i.e. $\frac{dx^\mu}{d\tau} = 0$.

As we know, we have to use covariant derivative for curved space, let us take the covariant derivative of this tangent vector

$$\begin{aligned} \nabla \frac{dx^\mu}{d\tau} &= \frac{\partial}{\partial \tau} \frac{dx^\mu}{d\tau} + \Gamma = 0 \\ \frac{\partial^2 x^\mu}{\partial \tau^2} &= -\Gamma = \text{acceleration} \end{aligned}$$

From Newton's law, we know $a = \frac{F}{m} = -\Gamma$. Now we know the Cristoffel symbol

$$\Gamma_{bc}^a(x) = \frac{1}{2} g^{ad} \left\{ \frac{\partial g_{dc}}{\partial x^b} + \frac{\partial g_{ab}}{\partial x^c} - \frac{\partial g_{bc}}{\partial x^d} \right\} \quad (9)$$

The general relativity and Newton's law of gravitation must coincide for ordinary masses, gravity and speed. When that happens, g^{ad} becomes 1 and the derivatives becomes very very small, near to zero except the time component $\frac{\partial g_{00}}{\partial x}$ of the metric tensor.

So above eqn. becomes

$$\Gamma_{bc}^a(x) = \frac{1}{2} \frac{\partial g_{00}}{\partial x} \equiv F$$

For unit mass

But in Newtonian mechanics, force is negative gradient of potential, i.e.,

$$F = -\frac{d\varphi}{dx} \quad \varphi = mgx$$

-ve sign indicates opposite direction of F and x

So the Cristoffel symbol

$$\Gamma_{bc}^a(x) = \frac{1}{2} \frac{\partial g_{00}}{\partial x} = \frac{d\varphi}{dx}$$

$$g_{00} = 2\varphi + \text{constant} \quad (9b)$$

Now the force capability across the whole sphere is

$$\int F \cdot dA = -\frac{GM}{r^2} \int dA = -\frac{GM}{r^2} 4\pi r^2 = -4\pi GM$$

Now from divergence theorem we can write,

$$\int \nabla \cdot \mathbf{F} dV = -4\pi G \int \rho dV$$

$$\nabla \cdot \mathbf{F} = -4\pi G\rho$$

$$\nabla^2 \varphi = 4\pi G\rho$$

Earlier we shown that $g_{00} = 2\varphi + \text{constant}$. Let we drop the constant, then,

$$g_{00} = 2\varphi$$

$$\varphi = \frac{1}{2} g_{00}$$

So we get,

$$\nabla^2 \frac{1}{2} g_{00} = 4\pi G\rho$$

$$\nabla^2 g_{00} = 8\pi G\rho \quad (9c)$$

Now the problem is, it is not a tensor equation, and for general relativity we need tensor equations. So what we looking for is something similar to it. So we want something that has the form $G_{\mu\nu}$, the Einstein's tensor, on the left side and $T_{\mu\nu}$, that contains all the mass, energy, pressure, stress terms, on the right side.

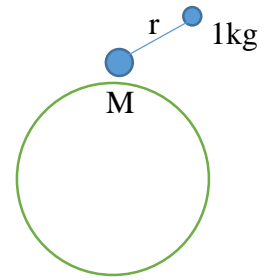
$$G_{\mu\nu} = 8\pi G T_{\mu\nu} \quad (9d)$$

This is essentially a vector but we need tensor $T_{\mu\nu}$

T_{00} is the time component of the tensor

T_{01}, T_{02}, T_{03} are called energy flow paths of the tensor

T_{00}	T_{01}	T_{02}	T_{03}
T_{10}	T_{11}	T_{12}	T_{13}
T_{20}	T_{21}	T_{22}	T_{23}
T_{30}	T_{31}	T_{32}	T_{33}



T_{10} , T_{20} , T_{30} are called momentum density of the tensor, and rest 9 components are essentially momentum, flux, stress, pressure parts of the tensor

Now the right hand side of the eqn. (9d) is $8\pi GT_{\mu\nu}$ which is mass term, so left side must be curvature term. Einstein thought this can be Ricci curvature tensor.

$$R_{\mu\nu} = 8\pi GT_{\mu\nu}$$

But the problem is, energy neither be created nor be destroyed.

If we take the derivative

$$\partial T_{\mu\nu} = 0$$

$$\partial R_{\mu\nu} \neq 0$$

We must always use the covariant derivative. So, if we got the covariant derivative of energy side is zero

$$\nabla(8\pi GT_{\mu\nu}) = 0 \tag{10}$$

$$\nabla T_{\mu\nu} = 0$$

we need something on the left side whose covariant derivative will also be zero. Einstein found that the covariant derivative of Ricci tensor is not zero but

$$\nabla R_{\mu\nu} = \frac{1}{2} \nabla g_{\mu\nu} R$$

$$\nabla \left(R_{\mu\nu} - \frac{1}{2} g_{\mu\nu} R \right) = 0 \tag{11}$$

So from eqn. 10 and 11

$$\nabla \left(R_{\mu\nu} - \frac{1}{2} g_{\mu\nu} R \right) = \nabla(8\pi GT_{\mu\nu}) \tag{12}$$

$$R_{\mu\nu} - \frac{1}{2} g_{\mu\nu} R = 8\pi GT_{\mu\nu}$$

For four dimensional purposes we have to divide right hand side of this eqn by C^4

$$R_{\mu\nu} - \frac{1}{2} g_{\mu\nu} R = \frac{8\pi GT_{\mu\nu}}{C^4}$$

But Einstein realized he forgotten something, he remembered that, $\nabla g_{\mu\nu} = 0$. Therefore they could have include in eqn. 12 with any constant in front of it.

$$\nabla \left(R_{\mu\nu} - \frac{1}{2} g_{\mu\nu} R + \Lambda g_{\mu\nu} \right) = \nabla \left(\frac{8\pi GT_{\mu\nu}}{C^4} \right)$$

$$R_{\mu\nu} - \frac{1}{2} g_{\mu\nu} R + \Lambda g_{\mu\nu} = \frac{8\pi GT_{\mu\nu}}{C^4}$$

First two term constitute the Einstein tensor. And Λ is called the cosmological constant, of which Einstein first thought of. When he was trying to identify how he could describe space in terms of mathematics